

Physics 407-07
Green's Function solution to wave equation

We want to find the Green's function solution to the equation

$$\square \Psi(t, \mathbf{x}; t', \mathbf{x}') = \delta(t - t') \delta^3(\mathbf{x} - \mathbf{x}') \quad (1)$$

where \mathbf{x} represents the three coordinates x, y, z .

We can take the Fourier transform of both sides

$$\Psi(\omega, \mathbf{k}) = \int \Psi(t, \mathbf{x}; 0, 0) e^{i(\omega t - \mathbf{k} \cdot \mathbf{x})} dt d^3x \quad (2)$$

where \mathbf{k} represents k_x, k_y, k_z , to get

$$(-\omega^2 + \mathbf{k} \cdot \mathbf{k}) \Psi(\omega, \mathbf{k}) = 1 \quad (3)$$

or

$$\Psi(\omega, \mathbf{k}) = \frac{-1}{\omega^2 - \mathbf{k} \cdot \mathbf{k}} \quad (4)$$

The

$$\Psi(t, \mathbf{x}; t' \mathbf{x}') = \int \int \int \int \frac{-1}{\omega^2 - \mathbf{k} \cdot \mathbf{k}} e^{-i(\omega(t-t') - \mathbf{k} \cdot (\mathbf{x} - \mathbf{x}'))} \frac{1}{(2\pi)^4} d\omega d^3\mathbf{k} \quad (5)$$

Writing $\mathbf{k} \cdot (\mathbf{x} - \mathbf{x}') = k|\mathbf{x} - \mathbf{x}'| \cos(\theta)$, where $k^2 = k_x^2 + k_y^2 + k_z^2$, θ is the angle between the \mathbf{k} vector and the \mathbf{x} vector, ϕ is the additional azimuthal angle, and $d^3\mathbf{k} = k^2 dk d(\cos(\theta)) d\phi$ we have

$$\Psi(t, \mathbf{x}; t' \mathbf{x}') = \frac{-1}{(2\pi)^3} \int \int \int \frac{e^{-i(\omega(t-t') + k|\mathbf{x} - \mathbf{x}'| \cos(\theta))}}{\omega^2 - k^2} k^2 dk d(\cos(\theta)) d\omega \quad (6)$$

$$= \frac{-1}{(2\pi)^3} \int \int \frac{1}{\omega^2 - k^2} e^{-i\omega(t-t')} \frac{e^{ik|\mathbf{x} - \mathbf{x}'|} - e^{-ik|\mathbf{x} - \mathbf{x}'|}}{ik|x - x'|} k^2 dk d\omega \quad (7)$$

$$= \frac{-1}{i|\mathbf{x} - \mathbf{x}'|(2\pi)^3} \int \int \frac{1}{\omega^2 - k^2} e^{-i\omega(t-t')} (e^{ik|\mathbf{x} - \mathbf{x}'|} - e^{-ik|\mathbf{x} - \mathbf{x}'|}) k dk d\omega \quad (8)$$

We now do the integral over ω by doing a contour integral. We want the contour at infinity to contribute nothing, so $e^{i\omega(t-t')}$ must go to zero as ω assumes imaginary parts at infinity. This means that we must close the

contour to positive imaginary ω for $t - t' > 0$ and to negative imaginary parts for $t - t' < 0$. We also want the integral to be zero for $t < t'$, so that the influence of the source is to the future, not the past. This means we must take the contour such that if we enclose it to negative imaginary ω , it must not enclose any of the singularities, which means that we need to take the contour along the real axis so that it runs below both of the singularities.

Thus

$$\begin{aligned}\Psi(t, \mathbf{x}; t'; \mathbf{x}') &= \frac{-1}{i|\mathbf{x} - \mathbf{x}'|(2\pi)^3} \int_0^\infty (2\pi i) \left(\frac{e^{-ik(t-t')}}{2k} - \frac{e^{ik(t-t')}}{2k} \right) (e^{ik|\mathbf{x}-\mathbf{x}'|} - e^{-ik|\mathbf{x}-\mathbf{x}'|}) k dk \Theta(t - t') \\ &= \frac{-1}{|\mathbf{x} - \mathbf{x}'|2(2\pi)^2} \int_0^\infty (-e^{ik(t-t'-|\mathbf{x}-\mathbf{x}'|)} + e^{-ik(t-t'-|\mathbf{x}-\mathbf{x}'|)}) + (e^{ik(t-t'+|\mathbf{x}-\mathbf{x}'|)} - e^{-ik(t-t'+|\mathbf{x}-\mathbf{x}'|)}) dk \\ &= \frac{\Theta(t - t')}{|\mathbf{x} - \mathbf{x}'|2(2\pi)^2} \int_{-\infty}^\infty e^{ik(t-t'-|\mathbf{x}-\mathbf{x}'|)} - e^{ik(t-t'+|\mathbf{x}-\mathbf{x}'|)} dk \\ &= \frac{\Theta(t - t')}{4\pi|\mathbf{x} - \mathbf{x}'|} (\delta(t - t' - |\mathbf{x} - \mathbf{x}'|) - \delta(t - t' + |\mathbf{x} - \mathbf{x}'|))\end{aligned}\tag{11}$$

where $\Theta(t - t') = 1$ if $t - t' > 0$ and is zero if $t - t' < 0$

But, since both $t - t'$ and $|\mathbf{x} - \mathbf{x}'|$ are positive, the second delta function can never have zero argument. This means that it is always zero. Thus

$$\Psi(t, \mathbf{x}; t'; \mathbf{x}') = \frac{\Theta(t - t')}{4\pi|\mathbf{x} - \mathbf{x}'|} \delta(t - t' - |\mathbf{x} - \mathbf{x}'|)\tag{13}$$

as required.

Note that there is a more invariant way of writing this

$$\Psi(t, \mathbf{x}; t'; \mathbf{x}') = -\frac{\Theta(t - t')}{2\pi} \delta((t - t')^2 - |\mathbf{x} - \mathbf{x}'|^2)\tag{14}$$

since

$$\int f(x) \delta(g(x)) dx = \int f(x(g)) \delta(g) \frac{dg}{dx} = \sum_i \frac{f(x_i)}{\frac{dg(x_i)}{dx}}\tag{15}$$

where $g(x_i) = 0$. Thus

$$\begin{aligned}&\int \frac{\Theta(t - t')}{2\pi} \delta((t - t')^2 - |x - x'|^2) F(t') dt' \\ &= \frac{1}{4\pi} \left(\frac{\Theta(|x - x'|) F(t - |x - x'|)}{|x - x'|} + \frac{\Theta(-|x - x'|) F(t + |x - x'|)}{-|x - x'|} \right)\end{aligned}\tag{16}$$

since $\partial_{t'}((t-t')^2 - |x-x'|^2) = -2(t-t')$ and evaluated at the two zeros of $((t-t')^2 - |x-x'|^2)$ for t' this is $-2(\pm|x-x'|)$

This latter form of the Green's function is far more clearly a Lorentz invariant form. The argument of the delta function is clearly Lorentz invariant, and the $\Theta(t-t')$ simply picks out the positive null cone rather than the negative, which is also clearly Lorentz invariant.